TA	BLE	I.	Various	equations	of	state.
----	-----	----	---------	-----------	----	--------

Equation	Acronym	Form
First-order Murnaghan	ME1	$z = \eta^{-1} \left(x^{\eta} - 1 \right)$
Second-order	ME ₂	$z = 2 (x^{(\eta^2 - 2\psi)^{1/2}} - 1)/$
Murnaghan	$(\eta^2 \ge 2\psi)$	$ \times [(\eta^2 - 2\psi)^{1/2} (\chi^{(\eta^2 - 2\psi)^{1/2} + 1}) - \eta (\chi^{(\eta^2 - 2\psi)^{1/2}} - 1)] $
Keane	KE	$z = [\eta^3 / (\eta^2 + \psi)^2]$
	$(-\eta^2 < \psi < 0)$	$\times (x^{(\eta^2 + \psi)/\eta} - 1) - [\psi/(\eta^2 + \psi)] \ln x$
First-order	BE1	$z = \frac{3}{2} \left(x^{7/3} - x^{5/3} \right)$
Birch		$\times [1 + \frac{3}{4} (\eta - 4) (x^{2/3} - 1)]$
Second-order	BE2	$z = \frac{3}{2} [x^{7/3} - x^{5/3}]$
Birch		$\times \{1 + \frac{3}{4}(\eta - 4) (x^{2/3} - 1)\}$
		$\frac{1}{24} [143 + 9\eta (\eta - 7) + 9\psi] \times (x^{2/3} - 1)^2]$
Grover, Getting, Kennedy	GGKE	$B_T = B_0^T \exp[\eta (1 - x^{-1})]$

two-parameter equations, the standard deviation of the V/V_0 -vs-P fit is less for the three-parameter equations. The weighted averages of the parameters B_0 , B'_0 , and B''_0 as determined by the two- and three-parameter equations of state are listed in Table IV. Also included in Table IV are the results of previous experiments for the purpose of comparison.

IV. DISCUSSION

A. Comparison with previous results

As can be seen from Table IV, present measurements of B_0 and B'_0 are not in agreement with those obtained earlier from static-compression techniques. The piston displacement method used by Vaidya and Kennedy is subject to errors caused by the finite yield strength of the NaCl, which is assumed to be absent and zero.

TABLE II. Bulk modulus and its pressure derivatives at 29.5°C.

Equation used ^a	B ₀ (kbar)	B ₀ '	B ₀ " (kbar ⁻¹)	Run No.
MEi	$238.14 \pm 0.10 \\ 238.20 \pm 0.12$	5.57 ± 0.04 5.60 ± 0.05	0.00 ^b 0.00	1 2
ME ₂	$237.72 \pm 0.24 \\ 237.89 \pm 0.20$	5.95 ± 0.20 5.96 ± 0.20	-0.11 ± 0.10 -0.13 ± 0.15	1 2
GGKE	$238.07 \pm 0.08 \\ 238.18 \pm 0.11$	5.64 ± 0.04 5.66 ± 0.05	-0.024 ^b -0.024 ^b	1 2
BE1	$238.04 \pm 0.08 \\ 238.15 \pm 0.11$	$5.68 \pm 0.04 \\ 5.69 \pm 0.05$	-0.035 ^b -0.035 ^b	1 2
BE2	$237.72 \pm 0.24 \\ 237.90 \pm 0.22$	5.97 ± 0.22 5.97 ± 0.22	-0.13 ± 0.10 -0.14 ± 0.15	1 2
KE	237.71 ± 0.24 237.89 ± 0.22	5.98 ± 0.20 5.98 ± 0.24	-0.14 ± 0.10 -0.15 ± 0.15	1 2

^aAcronyms defined in Table I.

^bObtained from B_0 and B'_0 using the appropriate expression given in Ref. 27.

TABLE III. Bulk modulus and its pressure derivatives at 40.4 °C.

Equation used ^a	B ₀ (kbar)	B ₀ '	<i>B</i> ₀ " (kbar ⁻¹)	Run No.
ME ₁	236.68 ± 0.08	5.53 ± 0.03 5.51 ± 0.03	0.00 ^b	1
	236.43 ± 0.32	5.55 ± 0.112	0.00 ^b	3
ME ₂	236.56 ± 0.22	5.65 ± 0.18	-0.04 ± 0.06	1
	236.28 ± 0.15 236.18 ± 0.41	5.85 ± 0.13 5.78 ± 0.35	-0.10 ± 0.15 -0.07 ± 0.20	23
GGKE	236.61 ± 0.08	5.61 ± 0.03	-0.024 ^b	1
	236.37 ± 0.33	5.62 ± 0.12	-0.024^{b}	3
BE1	236.59 ± 0.08	5.63 ± 0.03	-0.035 ^b	1
	236.56 ± 0.08 236.33 ± 0.34	5.61 ± 0.03 5.65 ± 0.12	-0.035^{b} -0.035^{b}	2
BE2	236.58 ± 0.22	5.64 ± 0.20	-0.04 ± 0.07	1
	236.28 ± 0.16 236.19 ± 0.43	5.87 ± 0.14 5.78 ± 0.38	-0.12 ± 0.15 -0.08 ± 0.24	23
KE	236.58 ± 0.22	5.64 ± 0.20	-0.04 ± 0.07	1
jo ej	$236.27 \pm 0.16 \\ 236.18 \pm 0.44$	$5.88 \pm 0.14 \\ 5.79 \pm 0.39$	$-0.13 \pm 0.15 \\ -0.08 \pm 0.25$	2 3

^aAcronyms defined in Table I.

^bObtained from B_0 and B_0' using the appropriate expression given in Ref. 27.

Singh and Kennedy²⁸ on the basis of x-ray studies suggest that the yield stress is quite high so that sizable deviatoric stresses are present. Ruoff²⁹ gives an estimate of the yield stress much smaller than Singh and Kennedy; nevertheless, even the presence of this yield stress would cause a sizable variation in the measured values of B_0 and particularly B'_0 using the piston displacement method. However, Kinsland and Bassett³⁰ in their x-ray studies do not observe a finite yield stress. It is conceivable that in the very long exposure time used by them considerable thermally activated stress relaxation occurred.

The agreement is fairly good when compared with ultrasonic measurements. The values obtained from B'_0 ultrasonically are generally lower than the present values, although they would tend to agree within the experimental uncertainties. In most of the ultrasonic work, errors are not listed with the data. Ghafelehbashi and Koliwad⁷ note that their values for derivatives of the individual directly measured elastic constants are good to within 5%. Since B involves a sum of two measured elastic constants $(B = C_1 - \frac{4}{3}C_s)$, B'_0 does likewise. Hence the potential for error is even larger. Even if an error of only 2.5% is used for the individual measured derivatives, the error in B'_0 would be about ± 0.4 . Most of these ultrasonic measurements were carried out only to a maximum of 3 to 4 kbar. Generally the data were analyzed by ignoring the contribution of B_0'' (except for Spetzler *et al.*) and therefore they would generally tend to yield lower values of B'_0 ; this neglect, for $B_0'' = -0.1$ kbar⁻¹, would lead to an ultrasonic value of B'_0 too low by about 0.2. One can compute a value of -0.09 kbar⁻¹ for B_0'' from the data of Spetzler *et al.*,⁹ and our measurements give an average value of -0.10kbar⁻¹. Considering the large uncertainties involved, the agreement is rather astonishing.

TABLE IV. Comparison of isothermal bulk modulus and its pressure derivative with previous values.

Source	Technique	B ₀ (kbar)	B ₀	<i>B</i> ₀ " (kbar ⁻¹)
Present work ^a	Length	238.1	5.64	the state
Present work ^b	Length	237.7	5.97	-0.12
Present work c	Length	237.7	5.71	-0.10
Present work ^d	Length	237.8	5.85	-0.12
Haussühl	Ultrasonic	237.3		1.1.1.1
Ghafelehbashi and Koliwad	Ultrasonic	237.0	5.37	
Barsch and Chang	Ultrasonic	234.2	5.39	20
Spetzler et al.	Ultrasonic	238.0	5.35	-0.09
Bartels and Schuele	Ultrasonic	234.0	5.35	
Bridgman	Piston-volume	240.8	4.61	
Vaidya and Kennedy	Piston-volume	231.7	4.92	
Fritz et al.	Shock	237.3°	5.50	and in

^aA weighted average of the two-parameter equations at

 $T = 29.5 \,^{\circ}\text{C}$

^bA weighted average of the three-parameter equations at T = 29.5 °C.

°A weighted average for the parameters B'_0 from all equations and B''_0 from all the three-parameter equations at T=29.5and 40.4 °C.

^dA weighted averaged for the parameters B'_0 and B''_0 from Keane's equation at T=29.5 and 40.4 °C.

"This value was not measured by them; instead Haussühl's value was used.

B. Discussion of B_0''

The most interesting result of this experiment is the determination of $B_0^{"}$. The weighted average value for $B_0^{"}$ as indicated in Table IV is -0.10 ± 0.05 kbar⁻¹. Admittedly $B_0^{"}$ has a large error associated with it. The main source of error in determining $B_0^{"}$ is the pressure itself. As pointed out before, ²⁷ the magnitude of $B_0^{"}$ depends on the functional form of the variation of P with the change in resistance per unit resistance ($\Delta R/R_0$) of the manganin gauge. As shown in Table V if one makes the assumption that the nonlinear pressure variation with the change in resistance of the manganin gauge is cubic, the weighted average value for $B_0^{"}$ is then -0.03 kbar⁻¹. Table V is included in the text to emphasize the sensitivity of $B_0^{"}$ to the possible uncertainty in pressure in this low-pressure region. What evidence is

TABLE V. Isothermal bulk modulus and its pressure derivative based on the assumption that the nonlinear pressure variation with the change in resistance of the manganin gauge is cubic.

B ₀ (kbar)	B ₀ '	B ₀ " (kbar ⁻¹)
238.20 ^a	5.60ª	The West Cana Trilling Diversity of
238.00 ^b	5.75 b	-0.06 ^b
238.00 °	5.53 °	-0.03°

^aA weighted average of the two-parameter equation at T = 29, 5 °C.

^bA weighted average of the three-parameter equations at T = 29.5 °C.

°A weighted average for the parameters B'_0 from all equations and B''_0 from all the three-parameter equations at T=29.5and 40.4°C. presently available (free piston data to 40 kbar)³¹ suggests that of three possible fits to the nonlinear term, quadratic, quadratic and cubic, and cubic, the quadratic fit is best. This does not rule out other important possibilities and must be considered an area where new developments and extra precision are needed. However, these results would indicate that B_0'' is more likely to be -0.10 kbar⁻¹. As expected, the dependence of B_0'' on the pressure variation of the manganin gauge is not as sensitive for low-bulk-modulus materials (e.g., sodium, potassium) as it is for high-bulk-modulus materials. The data for LiF indicates²⁷ even a much more sensitive dependence on the functional form of the gauge than NaCl does. B_0'' is related to the third pressure derivative of a function expressing the pressure in terms of the volume. Hence any error that is inherent in the pressure-vs-volume measurements will be propagated and compounded when B_0'' is computed. In other words the pressure-volume measurement must be extremely precise. The present V/V_0 measurements already have a precision of 1×10^{-7} . It is the pressure that we know only to a precision of 1×10^{-4} , as the mercury point³² is known only to 1 bar. Hence unless the pressure is measured to an extremely high precision of 1×10^{-6} , the error in B_0'' is going to remain at a relatively large magnitude.

C. Transition pressures of some fixed points

Using the values of B_0 , B'_0 , B''_0 that are listed in Table IV, the pressure transitions for the following transformations are estimated on the basis of two- and threeparameter equations of state and are tabulated in Table VI: (i) the barium I-II transformation, (ii) the bismuth III-V transformation, (iii) the transformation of bccphase iron to hcp phase, (iv) the transformation of NaCl itself from the NaCl to the CsCl structure.

The experimental values of V/V_0 used to estimate the

TABLE VI. Transformation pressures for some transitions and its comparison with previous values.

Equation used ^a	Ba I-II	Bi III-V b	Fe transition (kbar)	NaCl transition
ME ₁ ^b	59.9	86.2	170	468
BE1b	58.0	81.7	153	362
GGKE b	58.3	82.4	154	356
ME ₂ °	52.1	67.5	96	120
BE ₂ °	54.9	74.4	121	173
KE c	55.8	76.8	134	271
ME ₂ ^d	52.3	68.7	100	132
BE ₂ d	54.6	74.3	122	196
KEd	55.4	76.2	132	274
KE e	55.3	75.8	130	262
KE f	54.1	74.0	136	258
Decker	54.7	76.4	136	306
Drickamer		73-75	(110-113)	
Ref. 37	55	77	the standing to	1 the second

^aAcronyms defined in Table I.

^bThe values used for B_0 , B'_0 , and B''_0 are defined by footnote a in Table IV.

^cValues defined by footnote b in Table IV.

^dValues defined by footnote c in Table IV.

• Values defined by footnote d in Table IV.

^f Indicates that the values for B_0 , B_0' , and B_0'' obtained by Spetzler *et al.*, that are listed in Table IV, were used.